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# The GAPS programme with HARPS-N at TNG

## XI. Pr 0211 in M 44: the first multi-planet system in an open cluster <sup>★</sup>

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### ABSTRACT

**Context.** Open cluster (OC) stars share the same age and metallicity, and, in general, their age and mass can be estimated with higher precision than for field stars. For this reason, OCs are considered an important laboratory to study the relation between the physical properties of the planets and those of their host stars, and the evolution of planetary systems. However, only a handful of planets have been discovered around OC main-sequence stars so far, all of them in single-planet systems. For this reason we started an observational campaign within the GAPS collaboration to search for and characterize planets in OCs

**Aims.** We monitored the Praesepe member Pr 0211 to improve our knowledge of the eccentricity of the hot Jupiter (HJ) that is already known to orbit this star and search for additional intermediate-mass planets. An eccentric orbit for the HJ would support a planet-planet scattering process rather than a disk-driven migration after its formation.

**Methods.** From 2012 to 2015, we collected 70 radial velocity (RV) measurements with HARPS-N and 36 with TRES of Pr 0211. Simultaneous photometric observations were carried out with the robotic STELLA telescope to characterize the stellar activity. We discovered a long-term trend in the RV residuals that we show as being due to the presence of a second, massive, outer planet. Orbital parameters for the two planets are derived by simultaneously fitting RVs and photometric light curves, with the activity signal modelled as a series of sinusoids at the rotational period of the star and its harmonics.

**Results.** We confirm that Pr 0211b has a nearly circular orbit ( $e = 0.02 \pm 0.01$ ), with an improvement of a factor two with respect to the previous determination of its eccentricity, and estimate that Pr 0211c has a mass  $M_p \sin i = 7.9 \pm 0.2 M_J$ , a period  $P > 3500$  days and a very eccentric orbit ( $e > 0.60$ ). This kind of peculiar system may be typical of open clusters if the planet-planet scattering phase, which lead to the formation of HJs, is caused by stellar encounters rather than by unstable primordial orbits. Pr 0211 is the first multi-planet system discovered around an OC star.

**Key words.** (Stars:) individual: Pr 0211 — Stars: planetary systems — techniques: radial velocities, photometric

### 1. Introduction

Nearly 2 000 extrasolar planets in 1 300 planetary systems have been confirmed so far using a variety of techniques, prominently radial velocities (RV) and transit photometry (exoplanet.org, Han et al. 2014). Almost all these planets belong to isolated field stars, whose parameters (distance and age, above all) can be very uncertain. Consequently, the stellar mass and radius are affected

by large errors, which directly influence the precision with which we can estimate planet parameters (Casagrande et al. 2011; Santos et al. 2013). Comparing dynamical models of planet formation and evolution with observations requires a good knowledge of the age of the planetary systems, which in turn can be precisely determined only for a small sample of stars using asteroseismology. The chemical signature of planet formation on the stellar host is still an open question (see, e. g., González Hernández et al. 2013), with the interpretation of the results relying strongly on the choice of the stellar sample and the adopted methodologies, see, for example, the opposite points of view of Adibekyan et al. (2014) and Ramírez et al. (2014).

In principle, searching for planets in star clusters, in particular open clusters (OCs), offers an interesting and more appropriate laboratory. OC distances and ages can be precisely determined using statistical approaches, such as isochrone fitting and gyrochronology, which is more difficult to apply to field stars (Soderblom 2010). FGK main-sequence (MS) stars are the best

<sup>★</sup> Based on data obtained with (i) the HARPS-N spectrograph on the Italian Telescopio Nazionale Galileo (TNG), operated on the island of La Palma by the INAF - Fundacion Galileo Galilei (Spanish Observatory of Roque de los Muchachos of the IAC); (ii) the Tillinghast Reflector Echelle Spectrograph (TRES) on the 1.5-meter Tillinghast telescope, located at the Smithsonian Astrophysical Observatory's Fred L. Whipple Observatory on Mt. Hopkins in Arizona; (iii) the STELLA robotic telescopes in Tenerife, an AIP facility jointly operated by AIP and IAC. Figure 3, 4, 5, 9, 10 and 12, and Table 6 are made available in the online version of the paper. Table 1, 3, 4 and 5 are made available on-line at the CDS (<http://cdsweb.u-strasbg.fr/>).

targets, as they have low stellar activity, negligible mass loss, and make (low mass) planet identification easier, and planet-star interaction more straightforward to measure than in, for example, giant, evolved stars. Typically OC stars are chemically homogeneous so we can investigate the effect of the presence of planetary systems on the host star chemistry effectively. From the comparison of rotation and activity properties of planet hosts with those of other cluster members, we can get robust inference on the possible role of close-in planets on the evolution of stellar angular momentum (Lanza 2010).

Despite the many advantages offered by OCs, the number of planets discovered in OCs is still very small. An early RV search for planets in the Hyades cluster was attempted by Cochran et al. (2002) but the limited knowledge of the influence of stellar activity on RV measurements hampered the results. After the discovery of two long-period massive planets around OC giant stars by Sato et al. (2007) and Lovis & Mayor (2007), only recently new successes have been achieved with the discovery of two hot Jupiters (HJ) around two MS stars in M 44 (Quinn et al. 2012, hereafter Q12), two hot Neptunes around MS stars and a long-period giant planet around a sub-giant branch star in M 67 (Brucalassi et al. 2014), and an eccentric HJ around an MS star in the Hyades (Quinn et al. 2014). Except for the latter, eccentricities for all these planets have large uncertainties.

In this paper, we present the first results of an observing campaign within the Global Architecture of Planetary Systems (GAPS, Covino et al. 2013; Desidera et al. 2014; Poretti et al. 2015) for the search of exoplanets around 60 stars in the open clusters M 44 (Praesepe), Hyades and NGC 752. M 44 is located at 187 parsec ( $m - M = 6.36$  mag), with a colour excess of  $E(B - V) = 0.009$  mag and an age of  $790 \pm 30$  Myr that has been estimated from isochrone fitting (Brandt & Huang 2015). An independent age estimation of  $578 \pm 12$  Myr has been performed by Delorme et al. (2011) using gyrochronology. We focus on Pr 0211, a  $V = 12.15$ , G9V star hosting one of the two HJ discovered by Q12. We derive improved orbital parameters of Pr 0211b, a  $1.8 M_J$  planet with orbital period  $P = 2.14$  days and a RV semi-amplitude  $K \approx 300 \text{ m s}^{-1}$ . We also show that the same star hosts a second, more massive planet, located at a larger distance and in a much more eccentric orbit than Pr 0211b.

## 2. Observations

In 2012 we started an observational campaign with HARPS-N at the Telescopio Nazionale Galileo (TNG) targeting GK stars in the open cluster M 44, as part of GAPS. Our aims were to derive accurate and reliable eccentricities for the two known planets in the cluster, and search for additional intermediate-mass planets by pushing HARPS-N down to the limits imposed by stellar activity. Both goals can only be achieved with a careful observational strategy, because of the large activity of M 44 stars. During the observability season of the cluster, we took series of spectra on contiguous nights whenever the weather allowed us, and avoided observing the targets for just one or two isolated nights allocated to GAPS, to properly sample the rotational period of the stars ( $7.97 \pm 0.07$  days, Kovács et al. 2014, hereafter K14).

We started observing Pr 0211 in March 2013 and we continued observing the star until May 2015. It soon became clear that a long-term trend was present in the HARPS-N data, and therefore we further coordinated our RV campaign with TRES observations to reliably measure the RV offset between the instruments and take advantage of the earlier (2012) observations of Q12. A nearly simultaneous photometric campaign was conducted to characterize the stellar activity. The simultaneous char-

acterization of the eccentricity of the HJ, the quasi-periodic nature of stellar activity signal, and a possible long-period planet require an important investment of observing time spread across several years.

### 2.1. Radial velocity observations. I. HARPS-N dataset

We collected 70 HARPS-N spectra between March 2013 and May 2015. An exposure time of 1200 s was used to ensure an average signal-to-noise ratio (S/N) of 25 per extracted pixel at  $5500 \text{ \AA}$ , which resulted in an average radial velocity precision of  $6 \text{ m s}^{-1}$ . We did not make use of the simultaneous Thorium-Argon calibration lamp so as not to contaminate the stellar spectra, and because the RV noise for a star of this magnitude combined with the moderate rotation of the star ( $v \sin i_* = 4.8 \pm 0.5 \text{ km s}^{-1}$ , from Q12), largely dominates the internal stability of the instrument ( $\approx 1.0 \text{ m s}^{-1}$ , Cosentino et al. 2014). As a measure of precaution, we used tight constraints for the Moon position and the sky background illumination. The target was not observed when the distance from the Moon was less than  $60^\circ$  or when the sky was brighter than  $V = 19 \text{ mag arcsec}^{-2}$ .

Data reduction and RV extraction were performed using the most recent HARPS-N Data Reduction Software. We used the standard G2 mask to determine the weighted cross-correlation function (CCF, Baranne et al. 1996; Pepe et al. 2002). We used the CCF noise provided by the pipeline as estimate of the RV uncertainty.

### 2.2. Radial velocity observations. II. TRES dataset

The dataset included here is a continuation of the one presented in Q12, obtained with the Tillinghast Reflector Echelle Spectrograph (TRES, Fűrész 2008) between January and April 2012, one year before the HARPS-N observations. We refer to Q12 for a detailed description of the instrument and the data reduction details. In addition to the 18 RVs of the discovery paper, 36 new RVs with an average error of  $\approx 20 \text{ m s}^{-1}$  were gathered during the same temporal window of HARPS-N observations. TRES observations cover a time span of  $\approx 1200$  days i. e., 3.3 years.

### 2.3. Photometry

Concurrently with the RV campaign we performed a photometric follow-up to determine the rotational period of the star and monitor its long-term activity level. Observations were gathered with the *Wide Field STELLA Imaging Photometer* (WiFSIP) mounted on *STELLAR Activity-I* telescope at Teide Observatory, in Tenerife (Strassmeier et al. 2004; Weber et al. 2012). We obtained data spread over two months for every year between 2013 and 2015, with measurements in the Bessel  $B$  band over four to six separate pointings for a total of 56 (2013), 26 (2014), and 31 (2015) epochs. During 2015 the star was observed in the Sloan  $r$  band as well.

Data reduction, aperture photometry, and correction of light curves for systematics were performed following Nascimbeni et al. (2014). We used the generalized Lomb-Scargle periodogram (GLS, Zechmeister & Kürster 2009) to verify that the rotational period of  $\approx 8$  days was visible in our data.

## 3. Stellar parameters

Atmospheric parameters have been determined using the classical equivalent width (EW) approach on the HARPS-N coadded

spectrum with  $S/N \approx 320$ . EWs were measured using the latest version of ARES (Sousa et al. 2015) in automatic mode. Atmospheric parameters were determined using the 2014 version of the local thermodynamic equilibrium code MOOG<sup>1</sup> (Snedden 1973) and the Kurucz model atmospheric grid<sup>2</sup> with the new opacity distribution function (ODFNEW, Castelli & Kurucz 2004; Kurucz 1992). We used the iron linelist from Sousa et al. (2011), but we modified the values of the atomic line oscillator strength  $\log gf$  to match the solar EWs with the elemental abundances of Asplund et al. (2009) (Table 1).

Effective temperature  $T_{\text{eff}}$  and microturbulent velocity  $\xi$  were determined by minimizing the trend of iron abundances from individual lines with respect to their excitation potential and reduced EW respectively, while the surface gravity  $\log g$  was derived by imposing the same abundances for neutral and ionized iron lines. The final parameters with their errors were determined iteratively as in Dumusque et al. (2014). We checked the reliability of the automatic EW measurement by determining the photospheric parameters with the EWs extracted for a range of the  $\text{rejt}$  parameter, and verifying that the value automatically calculated by ARES was in fact providing the photospheric parameters with the lowest error. We found  $T_{\text{eff}} = 5270 \pm 60$  K,  $\log g = 4.46 \pm 0.10$ ,  $\xi = 1.1 \pm 0.1$  km s<sup>-1</sup> and  $[\text{Fe}/\text{H}] = 0.18 \pm 0.04$ , where the quoted uncertainties include possible systematic errors associated with this method (Sousa et al. 2008). These values are consistent within the errors with the ones listed by Q12 and obtained using a synthetic match approach (SPC, Buchhave et al. 2012), hence we used the weighted average of the two determinations as our final set of photospheric parameters in Table 2. We determined the projected rotational velocity  $v \sin i_{\star}$  by fitting synthetic spectra of seven isolated iron lines on the HARPS-N coadded spectrum of Pr 0211. We assumed the atmospheric parameters quoted in Table 2, a spectrograph resolution of  $R = 115000$ , a macroturbulent velocity  $v_{\text{mac}} = 2.3 \pm 0.7$  km s<sup>-1</sup> from the calibration of Doyle et al. (2014), and a limb darkening coefficient of 0.65 (Gray 2005). We obtain  $v \sin i_{\star} = 5.1 \pm 0.3$  km s<sup>-1</sup>, where the uncertainty includes the error contribution from  $v_{\text{mac}}$ .

The mass and radius of the star were derived using four sets of isochrones to analyse systematic differences between stellar models: PARSEC (Bressan et al. 2012), Dartmouth (Dotter et al. 2008), Yonsei-Yale (Yi et al. 2001) and BaSTI (Pietrinferni et al. 2004). We followed a Monte Carlo (MC) approach to take the uncertainties in temperature and metallicity into account. We took the average of the values as final parameters, while the associated errors are obtained by adding in quadrature the sample standard deviation of the measurements and their average error. We obtained a stellar mass of  $M_{\star} = 0.935 \pm 0.013$   $M_{\odot}$  and a radius of  $R_{\star} = 0.827 \pm 0.012$   $R_{\odot}$ .

We performed a second determination with the same approach using  $B - V = 0.87 \pm 0.01$ , which is an average of the values obtained by Johnson (1952) and Uggren et al. (1979) with a conservative choice for the associated errors, and the value for  $E(B - V)$  quoted in Section 1. In this case we obtained  $M_{\star} = 0.933 \pm 0.017$   $M_{\odot}$  and  $R_{\star} = 0.825 \pm 0.016$   $R_{\odot}$ , which was perfectly consistent with the spectroscopy-based value. Assuming an age of  $\approx 800$  Myr instead of  $\approx 600$  Myr caused a decrease of  $0.002$   $M_{\odot}$  in the derived masses and an increase of  $0.002$   $R_{\odot}$  for the stellar radius, i. e., the error on the age is negligible for the mass determination. In the following, we used the spectroscopy-based value for the mass and radius of the star.

**Table 2.** Parameters of the star.

Parameter	Value	Unit
$T_{\text{eff}}$	$5300 \pm 30$	K
$\log g$	$4.51 \pm 0.05$	dex
$\xi_t$	$1.1 \pm 0.1$	km s <sup>-1</sup>
$[\text{Fe}/\text{H}]$	$0.18 \pm 0.02$	dex
$v \sin i_{\star}$	$5.1 \pm 0.3$	km s <sup>-1</sup>
$M_{\star}$	$0.935 \pm 0.013$	$M_{\odot}$
$R_{\star}$	$0.827 \pm 0.012$	$R_{\odot}$
Age	$578 \pm 12^a$	Myr
	$790 \pm 30^b$	Myr
$i_{\star}$	$76 \pm 11$	deg
$\log R'_{\text{HK}}$	$-4.36 \pm 0.04$	dex

**References.** <sup>(a)</sup> Delorme et al. (2011). <sup>(b)</sup> Brandt & Huang (2015).

We took advantage of the knowledge of the photometric period  $P_{\text{rot}}$  from K14 and the projected rotational velocity  $v \sin i_{\star}$  to estimate the value of the inclination of the stellar spin axis  $i_{\star}$  to the line of sight. We included all the source of uncertainties through an MC approach, although we verified that the  $v \sin i_{\star}$  is the biggest contributor to our error budget.

#### 4. Analysis of RVs

The radial velocity data we collected from HARPS-N and TRES are shown in the upper panel of Figure 1, after removing the instrumental offset between the two datasets, as described in Section 6. The RVs, along with the activity indexes and photometry, are available online in the format shown in Table 3, 4, and 5. Before performing a full orbital fit we analysed the HARPS-N and TRES data independently with the GLS algorithm (second and third panel of Figure 1 respectively). The periodograms of both datasets confirm the presence of the planet with period of 2.14 days and RV semi-amplitude of  $\approx 300$  m s<sup>-1</sup>, as discovered by Q12. We also note that the alias at 1.86 days supplies much poorer folded RV curves.

The residuals obtained after subtracting a sinusoid at 2.14 days reveal the presence of a trend with the same amplitude in both datasets. The deviation of the trend from a sinusoidal shape makes the period estimate from GLS highly unreliable and a Keplerian curve should be used instead. However, before proceeding to an orbital fit, we need to assess if the long-period trend may be caused by stellar activity.

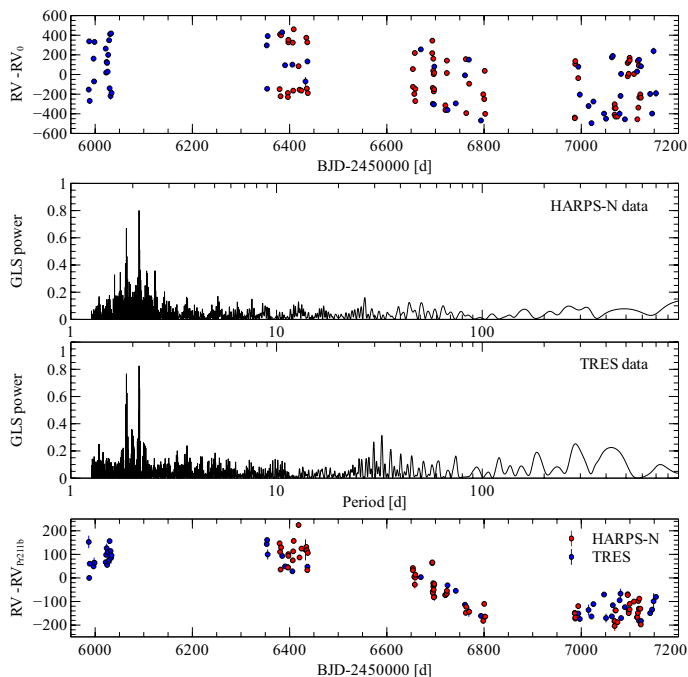
#### 5. Asymmetry and activity indexes

Line profile variation indicators were derived from the CCF using a variety of methods. The full width at half maximum (FWHM) of the CCF, the contrast of the CCF, and the bisector inverse slope (BIS) were determined automatically by the HARPS-N pipeline (Cosentino et al. 2012). The velocity span  $V_{\text{span}}$  (Boisse et al. 2011), the BIS<sup>+</sup> and BIS<sup>-</sup> (Figueira et al. 2013), and  $\Delta V$  (Nardetto et al. 2006) are derived using the scripts provided by Santos et al. (2014). We associated as conservative errors twice the  $\text{CCF}_{\text{noise}}$  for all these indicators. Additionally we used a modified version of the  $V_{\text{asy}}$  activity indicator (Lanza et al. in prep.<sup>3</sup>), which avoids the explicit presence of the RV in its def-

<sup>3</sup> The algorithm was introduced at the Extreme Precision Radial Velocity workshop and a description is available at <https://sites.google.com/a/yale.edu/eprv-posters/home>

<sup>1</sup> Available at <http://www.as.utexas.edu/~chris/moog.html>

<sup>2</sup> Available at <http://kurucz.harvard.edu/grids.html>



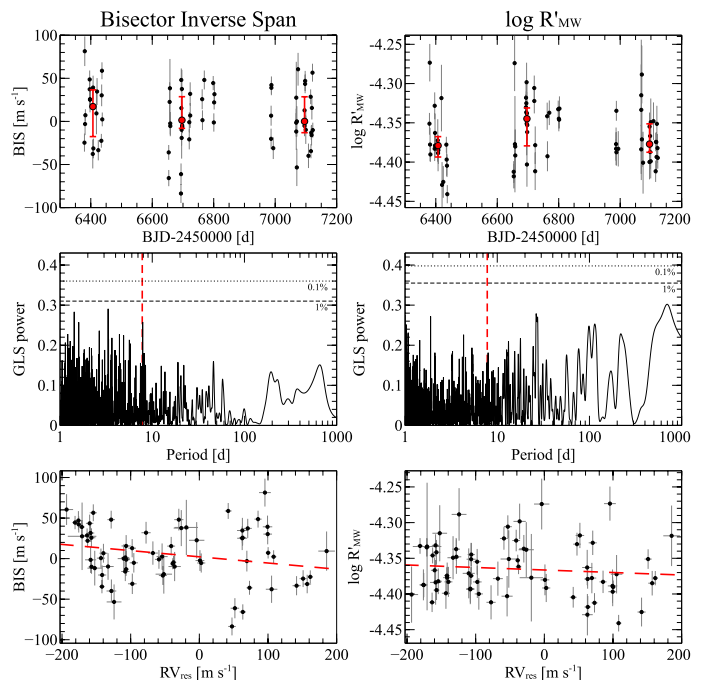
**Fig. 1.** Top panel: RV data obtained with HARPS-N (red dots) and TRES (blue dots) as a function of time. Middle panels: GLS periodogram of the observed RVs (black lines) for the HARPS-N and TRES datasets, respectively. Lower panel: residuals of the RVs after removing the signal from Pr 0211b using the solution found by GLS.

initiation, in contrast with the original formulation of Figueira et al. (2013).

In the analysis, we included activity indexes that are directly derived from the observed spectrum, rather than the CCF, and which are specifically designed to monitor the magnetic and chromospheric activity of the star. The  $\log R'_{\text{HK}}$  index (Noyes et al. 1984) was determined using a script that came with the HARPS-N DRS (Lovis et al. 2011), using the  $(B - V)_0$  color index quoted in Section 3. The  $H\alpha$  index (Gomes da Silva et al. 2011; Robertson et al. 2013), the NaID index (Díaz et al. 2007) and the HeI index (Boisse et al. 2009) were measured using our own implementation of the algorithms.

In Figure 2 we included only the BIS and  $\log R'_{\text{HK}}$  as representative of the other indexes listed above since they all show similar behavior. In the upper panels we plot the activity indexes as a function of time to check if there is any correlation with the long-period RV variation observed in the data (see the lower panels in Figure 1 for comparison). In the middle panels the GLS periodograms of the corresponding index are computed. The 1% and 0.1% false alarm probability (FAP) were estimated with a bootstrap approach.

In the lower panels of Figure 2 we investigated the presence of correlations between the asymmetry and activity indexes with RVs, after removing the signal of Pr 0211b. We determined the Spearman's rank correlation coefficient  $\rho$ , the slope of the linear fit  $m$  with its error and the  $p$ -value using the weighted least-square regression that was implemented in the StatsModels package<sup>4</sup>. Despite the signature of activity added to the RVs on the timescale of the rotational period, in all cases the correlation rank is close to zero, and the slope of the best fit is consistent with the lack of any trend within one or two  $\sigma$ , depending on the index that is considered (Table 6 and Figure 3).



**Fig. 2.** Analysis of the bisector inverse span (panels on the left side) and the  $\log R'_{\text{HK}}$  index (panels on the right side). Upper panels: the indicators as a function of time, the seasonal medians with the first and third quartiles indicated in red. Middle panels: GLS periodograms of the indexes, the rotational period of the star is indicated with a red vertical line. The 1% and 0.1% FAP levels are displayed as dashed and dotted horizontal lines, respectively. Lower panels: indicators as a function of RV, after removing the signal of Pr 0211b. The best fit is represented by the dashed red line

We assume here that any physical effect, whether spots or flares or magnetic cycles, capable of producing a peak-to-peak variation of  $\approx 280 \text{ m s}^{-1}$  in the RVs, should have left its imprint on at least one of the several indicators we have at our disposal. However, in all cases we can see that all the indicators are dominated by short-term variations, the effect of which has been averaged out by our observational strategy, see Section 2. While peaks at  $\approx 700$  days can be seen in the  $\Delta V$ ,  $\log R'_{\text{HK}}$  and  $H\alpha$  indexes, their significance is below the 1% threshold and they are not correlated to the longer period variation that is observed in the RVs (Figure 4 and 5). We note that the signal associated with the rotational period is visible in the periodograms, although not highly significant. This is, however, expected, given the low S/N of our measurements, combined with the slight loss of coherence of the short-period activity signal across several observing seasons, as already observed in Santos et al. (2014).

The marginal slope in the activity indexes advocates against a stellar origin for the long-term RV variation trend. While long-term trends in the chromospheric activity have been observed in several stars, their effect on RVs is usually well within  $100 \text{ m s}^{-1}$  and correlations between RV and activity indexes are seen even for moderately active stars (Santos et al. 2000). Despite the high level of activity of the star, the presence of a planet remains the simplest explanation to the long-term RV variation that we observed.

## 6. A two-planet system revealed

We performed the combined TRES and HARPS-N RV fit with the affine-invariant Markov Chain Monte Carlo (MCMC) en-

<sup>4</sup> Available at <http://statsmodels.sourceforge.net/>

semble sampler `emcee` code (Foreman-Mackey et al. 2013, hereafter FM13), coupled with the differential evolution genetic algorithm `PyDE`<sup>5</sup> to determine the global optimum solution as the initial point of the MCMC chains.

We take instrumental systematics into account by including an offset term between dataset, and adding in quadrature a jitter term to the errors that are associated with the measurements. Different jitter and offset parameters are used for datasets that come from different instruments, even if they are measuring the same quantity, and set as free parameters during the fitting process. We assumed the HARPS-N offset as the systemic velocity of the star  $\gamma$ .

We first analysed our data with a two-planet model that does not include any treatment for activity and works exclusively with the RVs that were obtained with TRES and HARPS-N. Our model has a total of 14 parameters, which are the Keplerian orbital parameters for two planets plus the systemic velocity of the star, a RV jitter term for each dataset, and finally the difference in the RV zero-point between HARPS-N and TRES data. We followed Eastman et al. (2013) by fitting the period  $P$  and the semi-amplitude  $K$  of the signal in the logarithmic space, and determining the eccentricity  $e$  and the argument of pericenter  $\omega$  by fitting  $\sqrt{e} \cos \omega$  and  $\sqrt{e} \sin \omega$ . Instead of using the time of periastron  $T_{\text{peri}}$ , which has to be constrained within one orbital period to facilitate chains convergence, we fitted the phase of the orbit  $\phi = M_0 + \omega$ , where  $M_0$  is the mean anomaly with respect to an arbitrary reference time  $T_{\text{ref}}$ .

Convergence of the chains is considered as achieved when the Gelman-Rubin statistic  $\hat{R}$  is lower than 1.03 (Gelman & Rubin 1992; Ford 2006). Additionally we verified that the acceptance fraction for all the walkers was in the range [0.20, 0.50], as suggested by FM13, as well as visual inspection of the individual chains. We discarded the section of the chains where the convergence criteria was not yet satisfied as burn-in phase. The median values and the standard deviation of the posterior distributions have been determined after applying a thinning factor of  $\simeq 100$  to the chains, i. e., the average autocorrelation time. We used the 34.13 percentile at each side of the median as an approximation of the standard deviation.

The results are listed in Table 7. For Pr 0211b, we determined a projected mass of  $1.88 \pm 0.03 M_J$  and an eccentricity that is consistent with zero. If we assume that the orbit of the planet and the spin of the star are aligned, as is expected for HJs around a star of  $T_{\text{eff}} \simeq 5300\text{K}$  (Dawson 2014), using our previous determination of  $i_*$  we obtain a true mass for the planet between 1.85 and  $2.05 M_J$ , well within the planetary mass range. For Pr 0211c, we determine a minimum mass of  $7.79 \pm 0.33 M_J$ , a very high eccentricity ( $e > 0.6$ ) and an orbital period in the range between 3000 and 9400 days. The effect of the activity can be envisaged in the large jitter term ( $\simeq 30 \text{ m s}^{-1}$ ) compared to the internal errors. This value is in agreement with the estimate provided by Santos et al. (2000) for the observed  $\log R'_{\text{HK}}$ . Activity is the cause of the prominent peak at 7.9 days in the periodograms of both HARPS-N and TRES RV residuals (Figure 6).

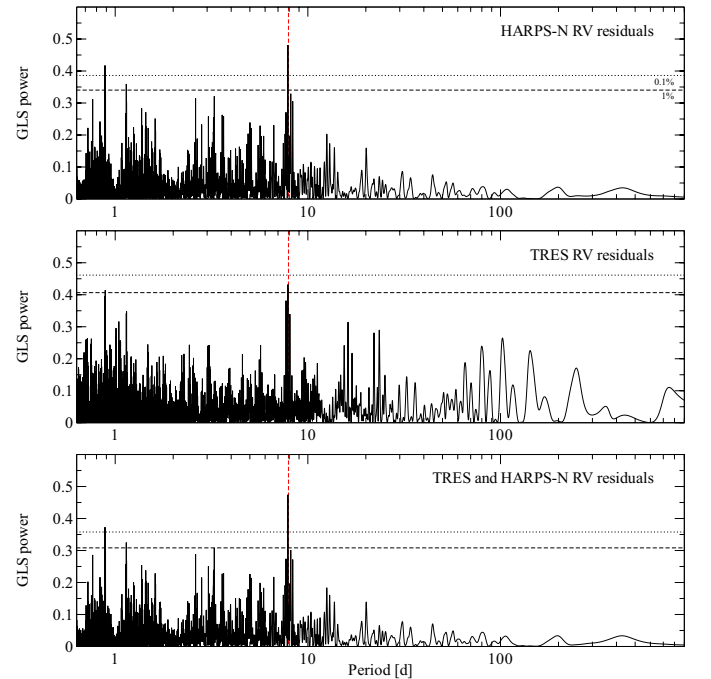
## 7. Simultaneous activity and orbital fit

In the previous section we included a RV jitter term to model both instrumental errors and stellar activity, thus implicitly assuming that the latter behaves as white noise. We know, however, that this assumption does not hold because of the correlated noise that was induced by activity variations, as seen in

**Table 7.** Orbital parameters of the two planets, obtained by fitting two Keplerian orbits and no signal for the activity. Best-fitting parameters are shown as fitted, additional parameters that were computed starting from the fitted parameters are tagged as derived.

Parameter	Pr 0211b	Pr 0211c	Note
P [days]	$2.14610 \pm 3 \cdot 10^{-5}$	$4850^{+4560}_{-1750}$	(a)
K [ $\text{m s}^{-1}$ ]	$309.7 \pm 4.2$	$138 \pm 7$	(a)
$\phi$ [deg]	$206.0 \pm 0.8$	$106.7 \pm 6.7$	(a)
$\sqrt{e} \sin \omega$	$0.016 \pm 0.083$	$0.78 \pm 0.09$	(a)
$\sqrt{e} \cos \omega$	$0.031^{+0.078}_{-0.086}$	$-0.30 \pm 0.11$	(a)
e	$0.011^{+0.012}_{-0.008}$	$0.71 \pm 0.11$	(b)
$\omega$ [deg]	$17^{+87}_{-111}$	$111 \pm 9$	(b)
$M_p \sin i$ [ $M_J$ ]	$1.88 \pm 0.03$	$7.79 \pm 0.33$	(b)
a [AU]	$0.03176 \pm 0.00015$	$5.5^{+3.0}_{-1.4}$	(b)
$T_{\text{peri}}$ [d]	$2456678.8 \pm 0.5$	$2456736 \pm 22$	(b)
Parameter	HARPS-N	TRES	kind
$\gamma$ [ $\text{m s}^{-1}$ ]	$35029 \pm 10$	$135 \pm 9$	(a)
RV <sub>jitter</sub> [ $\text{m s}^{-1}$ ]	$33 \pm 3$	$26 \pm 5$	(a)
$T_{\text{ref}}$	$2456679.97345$		(c)

**Notes.** (a) Fitted (b) derived (c) fixed.



**Fig. 6.** GLS periodograms of the RV residuals after subtracting the 2-planet solution, for HARPS-N data (upper panel), TRES data (middle panel) and the two datasets combined (lower panel). The rotational period of the star is indicated by a dotted vertical line.

photometry. In principle a better precision of the orbital parameters can be obtained if the activity is properly modeled. To do so, we developed a code<sup>6</sup> that can simultaneously handle datasets of different kinds, such as radial velocity, activity indexes, and photometry, as well as datasets that measure the same quantity but from different instruments, e. g., RVs obtained with two spectrographs or photometry obtained in several bands.

<sup>5</sup> Available at <https://github.com/hpparvi/PyDE>

<sup>6</sup> Available at <https://github.com/LucaMalavolta/PyORBIT>



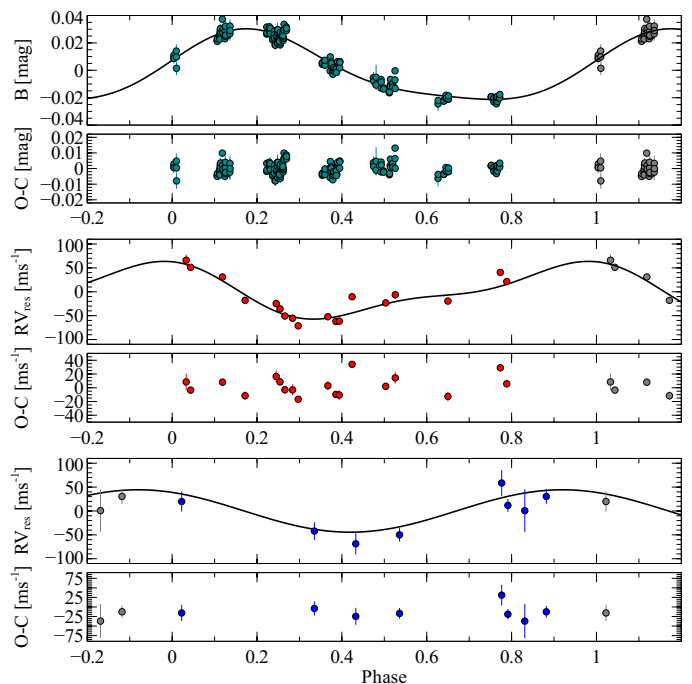
To include the effect of activity in our analysis, we followed the approach of Boisse et al. (2011). In this approach, the RV activity is modelled as the sum of sinusoidal signals at the rotational period  $P_{\text{rot}}$  of the star and its first harmonics. The rotational period and the phases of the sinusoids are constrained by the photometry, while the amplitude of the signal in RVs is left as a free parameter. The orbits of the planets and the activity signals are fitted simultaneously with the MCMC approach described in Section 6. In addition to the parameters previously introduced, we must include an offset term for each observational season in each photometric band, and two different jitters because of a substantial technical improvement of WiFSIP between the 2013 and 2014 observing seasons.

The rotational period and the phase of each sinusoid are common to all datasets, but we must take into account a possible delay on the effect of activity on the stellar line shape and the brightness of the star (Queloz et al. 2001). We simulated several spot configurations with SOAP 2.0 (Dumusque et al. 2014), using the astrophysical parameters of our star and the epochs of our datasets to sample the simulations, and verified that our activity model was able to reproduce the simulated data points.

The amplitudes of the sinusoids for HARPS-N data and TRES data are treated as independent parameters, since activity is a color-dependent effect and the two instruments are using different spectral ranges to determine the RVs. For the same reason, the amplitudes for photometric  $B$  and  $r$  bands are treated separately. Each season covered by photometric observations was fitted independently. Owing to the lack of simultaneous photometry, we excluded TRES data gathered in 2012 from the activity modeling and a separate RV jitter was introduced. Finally, for  $P_{\text{rot}}$ , we decided to use one free parameter for the whole temporal span since the periods for each observing season were consistent at  $1-\sigma$  within each other, when left as free parameters in the same fit. Each season is still characterized by independent values for the phases and amplitudes of the sinusoids. We assumed that the phases of the rotational signals are constant within an observational season, which is not strictly true, since these parameters are expected to change smoothly with time while spots/flares move on the surface of the star. We performed a preliminary fit using two harmonics of  $P_{\text{rot}}$ , then we excluded those harmonics which amplitudes were consistent with zero. We finally decided to use one harmonic for the HARPS-N dataset and just  $P_{\text{rot}}$  for the TRES RVs. In the fit we did not include any activity index, since no clear signal is visible at  $P_{\text{rot}}$  and its harmonics. As such, we have a total of 35 additional parameters. Results of the analysis are reported in Table 8.

Our treatment of the activity reduces the errors on the orbital parameters of the two planets up to  $\approx 30\%$ , as in the case of the semi-amplitudes of the RV signals. The improvement in the RV fit, with a strong reduction of the HARPS-N RV jitter (from  $33 \text{ m s}^{-1}$  before to  $15 \text{ m s}^{-1}$  after activity correction) and the TRES one (from  $26$  to  $13 \text{ m s}^{-1}$ ), is visually summarized in Figure 7, where the best-fit model and the RV residuals for the two planets, before and after taking activity into account, are compared side by side. The relationship between the photometric signal and the RVs (after removing the signal of the two planets) is shown in Figure 8, 9, and 10 for each observational season. The slight loss of coherence of the activity-induced signal in the asymmetry/activity indexes is supported by the change of shape and phase observed in the photometric light curves, as well as the different shapes of the RV sinusoids across the seasons.

The median values of the orbital parameters do not change significantly after including the activity model in the overall fit. This implies that the two observed planetary signals are not an



**Fig. 8.** Activity model for the data gathered in 2013. The upper panel shows the photometric light curve with residuals. Middle and lower panels show the RV data from HARPS-N and TRES, respectively (color coded as in Figure 1), after removing the two-planet solution. Data have been phased on the rotational period obtained from the global fit.

artifact that was introduced by our methodology to correct for the activity-induced RV variations. However, we note that this may not hold in general, and that Pr0211 represents a special case because of the large separation in frequency space between the planetary and activity signals.

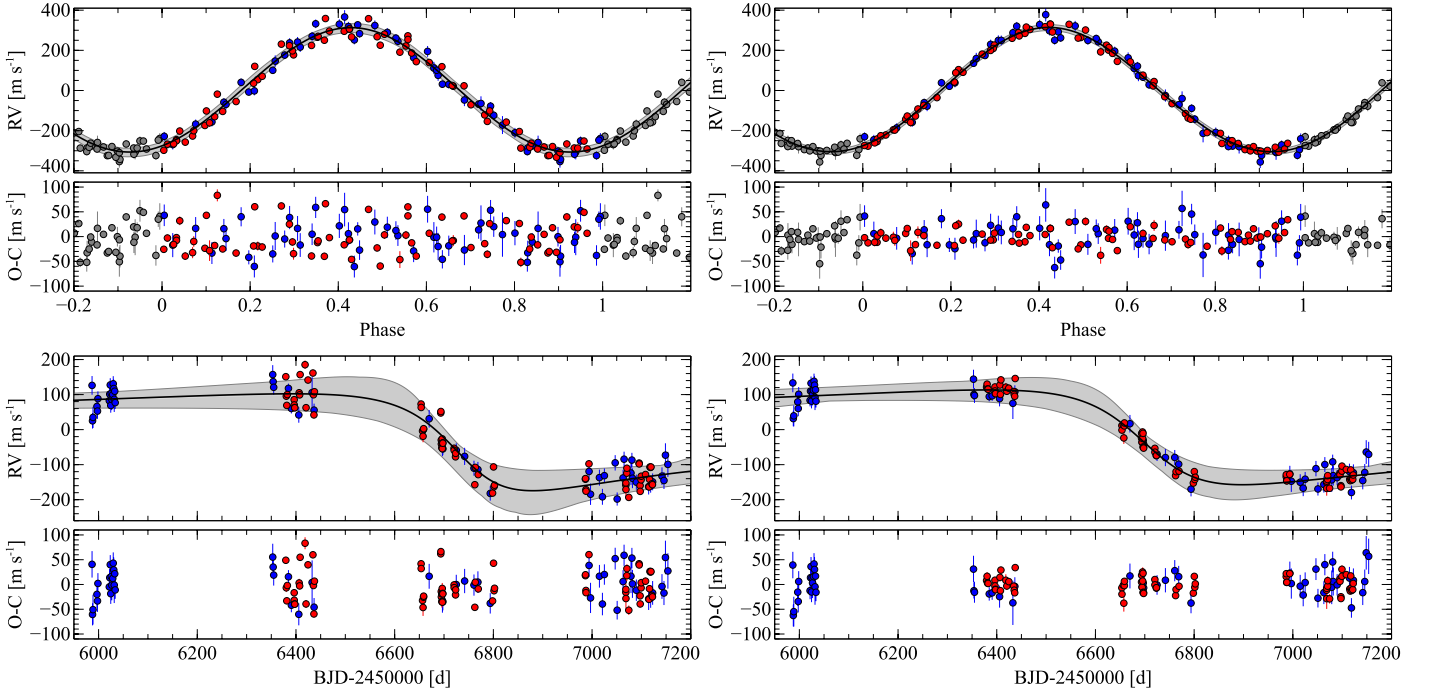
With an eccentricity of  $0.02 \pm 0.01$ , following the criterion of Lucy & Sweeney (1971), the inner planet can be safely considered as being on a circular orbit. Our determination improves the previous determination of Q12 by a factor of two ( $e = 0.05 \pm 0.02$ ).

The limited time-span of our observation allows us to only put a limit on the period of the outer between 9.6 and 27 years, and between 0.6 and 0.8 on its eccentricity, with a strong degeneracy between the two parameters. The projected mass range is well constrained at  $\approx 7.95 \pm 0.25 \text{ M}_J$  thanks to the observation of both the turning points in the RV curve. The relationship between these parameters is visualized in the density plot of the posterior distribution for Pr0211c, in Figure 11.

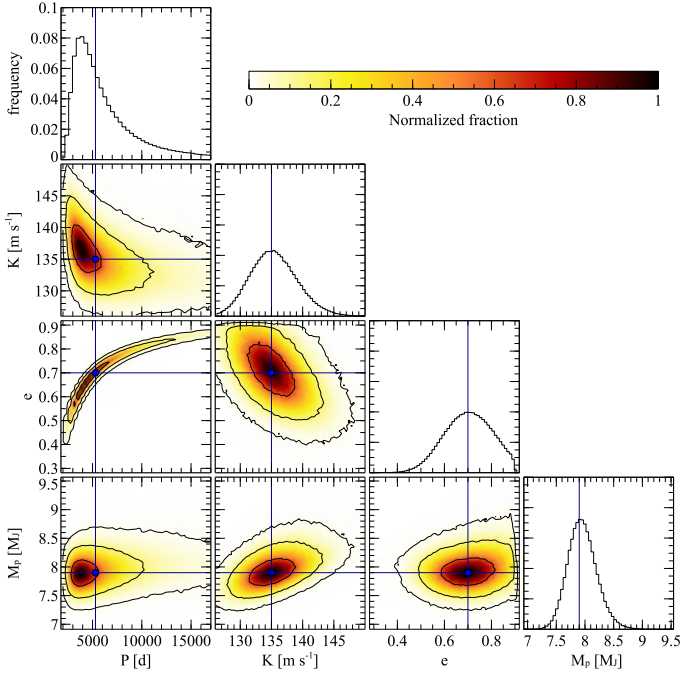
RV residuals do not show any significant peak in their periodogram (Figure 12). No additional signals that could be interpreted as planets are detected in our data. We can safely rule out the presence of a planet with an RV semi-amplitude greater than the observed RV jitter.

## 8. Discussion

The activity of Pr0211 is noticeable if compared to the average RV jitter of  $13 \text{ m s}^{-1}$  found by Q12 for this cluster, and  $16 \text{ m s}^{-1}$  found by Paulson et al. (2004) for the almost coeval Hyades cluster. For this star, photometric sinusoidal variations correspond to RV amplitudes up to  $50 \text{ m s}^{-1}$ , and the RV jitter term in the combined fit of the planet motion and activity signal for the HARPS-N RVs is still very high ( $\approx 15 \text{ m s}^{-1}$ ), compared to the average internal error of the RVs ( $\approx 6 \text{ m s}^{-1}$ ). One possibility could be



**Fig. 7.** Orbital solution and RV residuals before (left panels) and after (right panels) the inclusion of activity in the global fit. The two upper panels show the RV fit and its residuals for the inner planet, phased on its period, after removing the solution of the second planet. The two lower panels show the RV fit as a function of time for the outer planet, after removing the solution of the inner one. Red and blue dots represent HARPS-N and TRES data, respectively. The gray shaded areas represent the  $3\sigma$  confidence regions



**Fig. 11.** Posterior distribution of period  $P$ , RV semi-amplitude  $K$ , orbital phase, eccentricity  $e$ , and planetary mass for Pr 0211c. Histograms share the same vertical scale. Blue crosses identify the median values reported in Table 8.

that the star is subjected to star-planet interactions caused either by the inner planet, which could explain the short-term variability, but the low S/N of the spectra does not allow us to draw any conclusions on the possible connection between the two planets and their host star. The photometric amplitude of  $\approx 0.025$  mag

that we measure in the B filter during the first two seasons is higher than the 0.017 mag measured in white light by K14. This is expected since at shorter wavelengths the contrast between the cool stellar spots and the hotter surrounding photosphere is larger.

The dynamical stability of the system has been verified by running a 100 Myr simulation with SyMBA (Duncan et al. 1998). We used the values listed in Table 8 as initial parameters for the planets, performing three separate analyses with the median values and their  $\pm\sigma$  variations for the orbital parameters of Pr 0211c, and the median parameters for Pr 0211b.

We also estimated the characteristic tidal timescales of the host star and the inner planet using the model of Leconte et al. (2010) and assuming the stellar parameters in Table 2, the planetary parameters in Table 8, and a radius of  $1.0 R_J$  for the planet. We assumed modified tidal quality factors  $Q'_s = 10^7$  for the star (Jackson et al. 2009), and  $Q'_p = 10^5$  for the planet, similar to the value measured for Jupiter (Lainey et al. 2009). The timescale for the circularization of the orbit is ruled by the dissipation inside the planet and it is much shorter ( $\sim 15$  Myr) than the age of the system, thus supporting our result about a circular orbit in Section 7. The timescales of tidal orbital decay and star spin alignment are ruled by the tidal dissipation inside the star and are proportional to  $Q'_{\text{star}}$ . With the adopted value and the present system parameters, they are of  $\approx 6$  and  $\approx 11$  Gyr respectively, thus significantly longer than the age of the star. Therefore, a possible primordial misalignment between the stellar spin and the orbital angular momentum could still be present.

While the analysis of Kepler candidates seems to suggest that HJs do not generally have small-size companions (Steffen et al. 2009), several examples of HJs with periods shorter than 10 days and outer massive companions with  $P > 100$  do exist. For example, the planetary systems around  $\nu$  And (Butler et al. 1997), HD 187123 (Wright et al. 2007), HAT-P-13 (Bakos



et al. 2009), HD 217107 and HIP 14810 (Wright et al. 2009), Kepler-424 (Endl et al. 2014), WASP-41 and WASP-47 (Neveu-VanMalle et al. 2016). Among these systems, only HAT-P-13 and HD 217107 show a similar architecture to Pr 0211, i.e., the outer planet has high eccentricity ( $e > 0.5$ ). A recent investigation by Knutson et al. (2014), which was aimed at finding outer companions to HJs, estimated a frequency of  $50 \pm 10$  % outer planets with mass  $1-13 M_J$  within  $1-10$  AU from their hosting star. In fact, the observed statistics on the number of HJs with outer planets is strongly affected by several biases, both because the host star is usually excluded from the RV surveys after the discovery of an HJ, and because of the difficulty of disentangling long-term magnetic cycles from planetary signals (e.g., Damasso et al. 2015).

The present orbital architecture of Pr 0211 may be typical of a planetary system with at least three planets, which has experienced a period of chaotic dynamics evolving into planet-planet scattering. Eventually, two planets were left on stable, inner orbits, while the other ones are ejected from the system on a hyperbolic trajectory. Subsequent tidal interaction with the host star circularize the orbit of the inner planet to the periastron distance, transforming it into a HJ, while the outer one survives on an eccentric and misaligned orbit (Weidenschilling & Marzari 1996; Chatterjee et al. 2008; Nagasawa et al. 2008).

In this scenario, two viable processes triggered the planet-planet scattering phase: the planets have been on dynamically unstable orbits ever since their formation, or the planetary system experienced a stellar encounter in the initial evolution of the cluster. In the first case, the fraction of planetary systems in clusters that underwent a planet-planet scattering process would be comparable to that of field stars or only slightly higher. In the second case, since close stellar flybys are frequent in the initial stages of cluster evolution, the typical orbital architecture of planets around stars in clusters is expected to be shaped by planet-planet scattering events more frequently than for field stars (Zakamska & Tremaine 2004; Malmberg & Davies 2009). As a consequence, we should have a large number of systems with a HJ and a second giant planet on an outer eccentric orbit with respect to the second case. Recent simulations have tried to tackle the problem by determining the expected frequency of such systems (Hao et al. 2013; Li & Adams 2015) or advocating a cluster origin for several field stars with planetary systems similar to that of Pr 0211 (Shara et al. 2016).

On the other hand, the majority of cluster stars disperse into the field within the first 10 Myr (e.g., Fall et al. 2009; Dukes & Krumholz 2012), so stellar encounters are more probable during the early planetary formation phase when the star is still embedded in its nest disc. In this case, the dynamical excitation of a stellar encounter is fully damped by the circumstellar disc and any trace of the flyby is erased (Marzari & Picogna 2013; Picogna & Marzari 2014). The probability of stellar encounters at later time, however, could be increased if the cluster is the outcome of two smaller clusters merging, as it has been suggested for M 44 (Holland et al. 2000).

More planetary systems in open clusters are needed to clarify the contribution of stellar flybys to the final orbital configurations of multi-planet systems that were originally formed in clusters.

## 9. Conclusion

We have presented the discovery of the first planetary system around an open cluster star. New HARPS-N and TRES radial velocity measurements of the planet-host star Pr 0211 in M 44

have led us to the detection of a long-period, massive Jupiter and to confirm the nearly-circular orbit of the already known inner planet. A deep analysis of the stellar activity has been performed to exclude a stellar cause for the additional RV signal. We then performed a joint modeling of the activity and the planetary signals to improve the precision of the orbital parameters of the two planets, with the inclusion of a photometric light curve that had been specifically gathered for this purpose.

The discovery of the first multi-planet system in an open cluster is the first result of a long-term search for planets in OCs, which is conducted within the GAPS collaboration, and which will be described in a forthcoming paper.

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**Table 8.** Orbital parameters of the two planets, obtained by including activity in the global fit.  $K$  denotes the semi-amplitude of the RV sinusoid at the given harmonic of the stellar orbital period,  $A$  the semi-amplitude of the photometric sinusoids,  $\phi$  the phases of the sinusoids (in common between RV and photometry). Only one parameter has been used for the jitter term in the  $B$  light curves, obtained during 2014 and 2015.

Parameter	Pr 0211b	Pr 0211c	Unit
P	$2.14609 \pm 2 \cdot 10^{-5}$	$5300^{+4450}_{-1800}$	days
K	$309.4 \pm 2.5$	$135 \pm 4$	$\text{m s}^{-1}$
$\phi$	$207.0 \pm 0.5$	$101 \pm 5$	deg
$\sqrt{e} \sin \omega$	$-0.06^{+0.07}_{-0.05}$	$0.81^{+0.06}_{-0.07}$	
$\sqrt{e} \cos \omega$	$0.10^{+0.04}_{-0.07}$	$-0.19^{+0.09}_{-0.08}$	
$e^a$	$0.017 \pm 0.010$	$0.70 \pm 0.10$	
$\omega^a$	$329 \pm 35$	$103 \pm 6$	deg
$M_p \sin i^a$	$1.88 \pm 0.02$	$7.95 \pm 0.25$	$M_J$
$a^a$	$0.03184 \pm 0.00015$	$5.8^{+2.9}_{-1.4}$	AU
$T_{\text{peri}}^a$	$2456678.6 \pm 0.2$	$2456709 \pm 16$	days
Parameter	HARPS-N	TRES	unit
$\gamma$	$35035 \pm 8$	$128 \pm 7$	$\text{m s}^{-1}$
jitter RV	$15 \pm 2$	$13 \pm 7$	$\text{m s}^{-1}$
jitter RV <sub>2013</sub>	-	$19 \pm 9$	$\text{m s}^{-1}$
$P_{\text{rot}}$	$7.93 \pm 0.01$		days
RV-Phot. $\phi$	$263 \pm 4$		deg
$T_0^b$	$2456679.97345$		days
2013 obs. season			
$K_{P_{\text{rot}}}$	$50.7 \pm 6.6$	$44.4 \pm 10.6$	$\text{m s}^{-1}$
$K_{P_{\text{rot}}/2}$	$19.2 \pm 5.4$	-	$\text{m s}^{-1}$
$A_{P_{\text{rot}}}$	$0.0251 \pm 0.0005$		$B$ mag
$A_{P_{\text{rot}}/2}$	$0.0054 \pm 0.0007$		$B$ mag
$\phi_{P_{\text{rot}}}$	$119 \pm 15$		deg
$\phi_{P_{\text{rot}}/2}$	$73 \pm 30$		deg
offset $B$	$-0.0098 \pm 0.0004$		$B$ mag
jitter $B$	$0.0033 \pm 0.0003$		$B$ mag
2014 obs. season			
$K_{P_{\text{rot}}}$	$27.8 \pm 5.1$	$15.6 \pm 10.3$	$\text{m s}^{-1}$
$K_{P_{\text{rot}}/2}$	$33.8 \pm 5.0$	-	$\text{m s}^{-1}$
$A_{P_{\text{rot}}}$	$0.0246 \pm 0.0008$		$B$ mag
$A_{P_{\text{rot}}/2}$	$0.0038 \pm 0.0008$		$B$ mag
$\phi_{P_{\text{rot}}}$	$215 \pm 7$		deg
$\phi_{P_{\text{rot}}/2}$	$322 \pm 10$		deg
offset $B$	$0.0044 \pm 0.0006$		$B$ mag
jitter $B_{2014,2015}$	$0.0050 \pm 0.0003$		$B$ mag
offset $B_{\text{short}}$	$0.0034 \pm 0.0010$		$B$ mag
jitter $B_{\text{short}}$	$0.0082 \pm 0.0008$		$B$ mag
2015 obs. season			
$K_{P_{\text{rot}}}$	$32.4 \pm 5.2$	$28.6 \pm 9.2$	$\text{m s}^{-1}$
$K_{P_{\text{rot}}/2}$	$10.4 \pm 5.1$	-	$\text{m s}^{-1}$
$A_{P_{\text{rot}}}$	$0.0124 \pm 0.0008$		$B$ mag
$A_{P_{\text{rot}}/2}$	$0.0050 \pm 0.0008$		$B$ mag
$A_{P_{\text{rot}}}$	$0.0092 \pm 0.0006$		$r$ mag
$A_{P_{\text{rot}}/2}$	$0.0035 \pm 0.0005$		$r$ mag
$\phi_{P_{\text{rot}}}$	$312 \pm 22$		deg
$\phi_{P_{\text{rot}}/2}$	$79 \pm 43$		deg
offset $B$	$0.0045 \pm 0.0006$		$B$ mag
jitter $B_{2014,2015}$	$0.0050 \pm 0.0003$		$B$ mag
offset $r$	$0.0033 \pm 0.0004$		$r$ mag
jitter $r$	$0.0032 \pm 0.0003$		$r$ mag

**Notes.** All the parameters have been fitted except: <sup>(a)</sup> derived <sup>(b)</sup> fitted .

**Table 1.** Linelist with the oscillator strength  $\log gf$  updated to reflect the new solar abundances in the 2014 version of MOOG. All the values except the  $\log gf$  are taken from <http://www.astro.up.pt/~sousasag/ares/>. The full table is available on-line, a portion is shown here for reference.

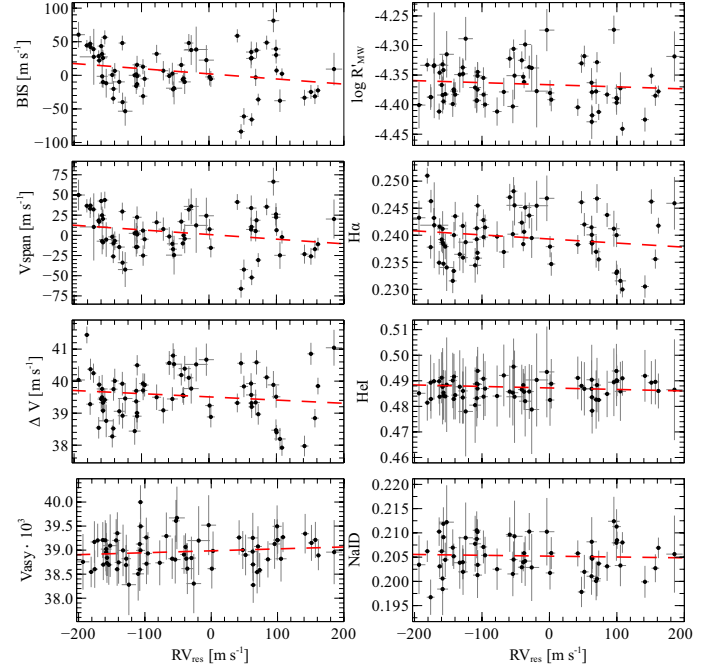
Element	$\lambda$ [Å]	$\chi$ [eV]	$\log gf$	EW [mÅ]	
FeII	26.1	4508.28	2.86	-2.539	87.3
FeII	26.1	4520.22	2.81	-2.688	81.9
FeI	26.0	4523.40	3.65	-1.958	44.2
...	...	...	...	...	...

**Table 3.** RV data collected with HARPS-N (a) and TRES (b). Instrumental offsets have not been applied. The full table is available on-line, a portion is shown here for reference.

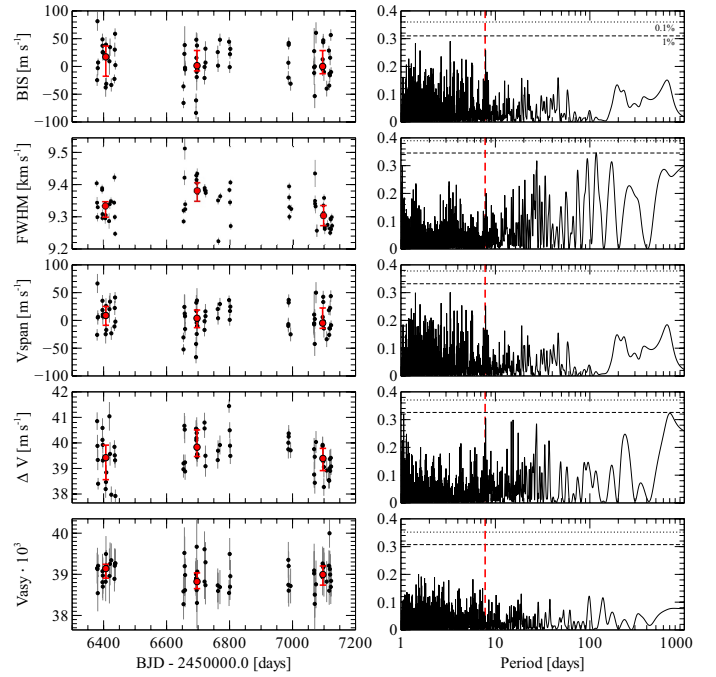
BJD <sub>UTC</sub> [d]	RV [m s <sup>-1</sup> ]	$\sigma$ RV [m s <sup>-1</sup> ]	Instrument
2456379.42879	34880.0	5.2	a
2456380.43289	35438.1	8.6	a
2456381.39909	34807.5	6.4	a
...	...	...	...

**Table 6.** Spearman’s correlation rank, slope of the linear fit and its  $p$ -value for several asymmetry and activity indexes as a function of RVs, after he signal of the first planet has been removed.

Index	$\rho$	$m$	$p$ -value
BIS	-0.23	$-0.08 \pm 0.04$	0.11
RV span	-0.21	$-0.05 \pm 0.03$	0.17
$\Delta V$	-0.18	$-0.001 \pm 0.001$	0.70
$V_{asy}$	0.17	$0.0004 \pm 0.001$	0.31
$\log R'_{HK}$	-0.10	$(-4 \pm 4) \cdot 10^{-5}$	0.39
H $\alpha$	-0.03	$(-7 \pm 5) \cdot 10^{-6}$	0.57
HeI	0.16	$(6 \pm 4) \cdot 10^{-6}$	0.19
NaID	-0.06	$(-2 \pm 4) \cdot 10^{-6}$	0.67



**Fig. 3.** CCF asymmetry indexes (panels on the left side) and chromospheric activity indicators (panels on the right side) as a function of RVs, after removing the signal of Pr 0211b. The best fit is represented by the dashed red line.



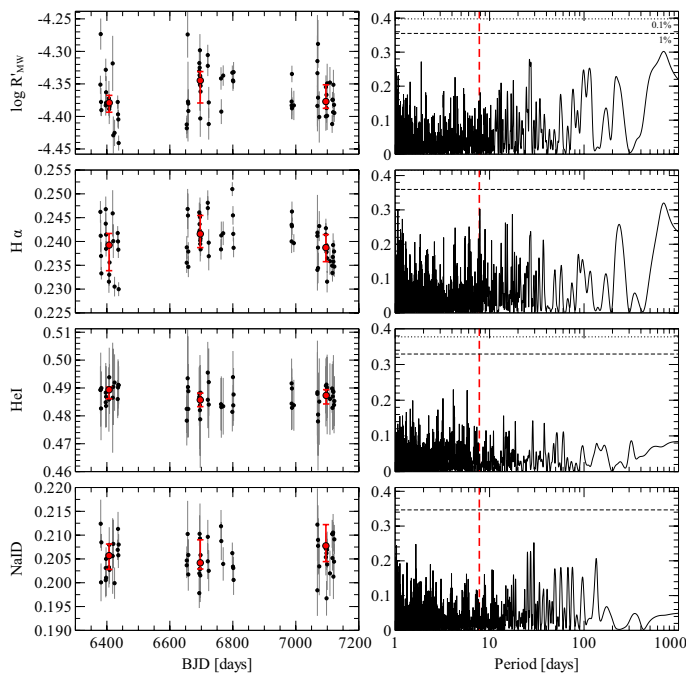
**Fig. 4.** Asymmetry indicators as a function of time (panels on the left) and their corresponding GLS periodogram (panels on the right). From top to bottom the bisector inverse span BIS, the FWHM of the CCF, the RV span, the  $\Delta V$  and  $V_{asy}$  are shown. Seasonal medians with the first and third quartiles are indicated in red. The rotational period of the star is indicated with a red vertical line. The 1% and 0.1% FAP levels are displayed as dashed and dotted horizontal lines, respectively. The extra power in the FWHM periodogram is due to a drift in the instrumental resolution, corrected in March 2014 (JD  $\approx$  2456370).

**Table 4.** CCF Asymmetry indicators extracted from HARPS-N spectra. The errors associated to BIS, BIS−, BIS+,  $\Delta V$  and  $V_{\text{span}}$  are twice  $\sigma_{\text{RV}}$ , so they have not been reported in the table. The full table is available online, only a portion is shown for references.

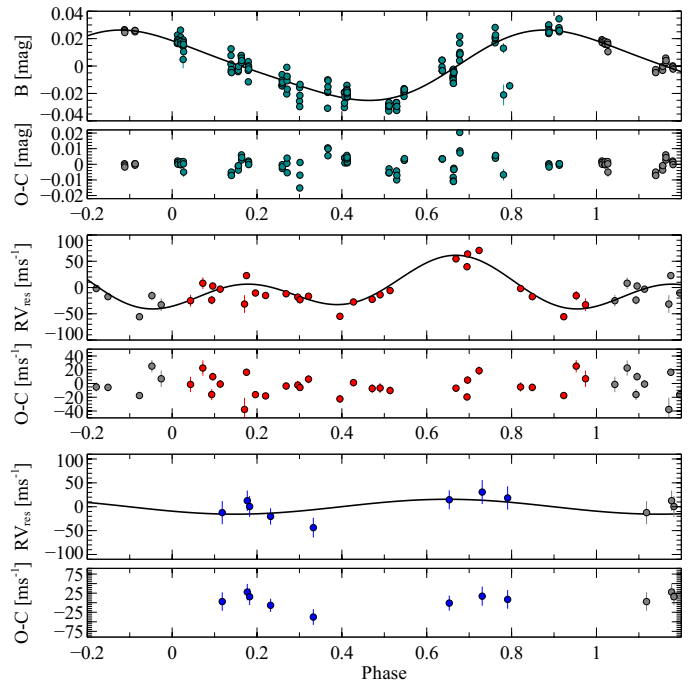
BJD-2450000.0	$\text{RV}_{\text{res}}$	$\sigma_{\text{RV}}$	BIS	BIS−	BIS+	$\Delta V$	$V_{\text{span}}$	FWHM	$\sigma_{\text{FWHM}}$	$V_{\text{asy}} \cdot 10^3$	$\sigma_{V_{\text{asy}}} \cdot 10^3$
6379.428774	151.14	5.2	−24.6	−13.3	−40.9	43.1	−26.0	9.403	0.010	39.15	0.37
6380.432878	95.54	8.6	81.4	52.6	86.2	−83.8	66.4	9.343	0.017	39.12	0.51
6381.399079	69.21	6.4	−3.0	2.3	−5.4	1.2	5.2	9.299	0.012	38.54	0.44
...	...	...	...	...	...	...	...	...	...	...	...

**Table 5.** Activity indicators extracted from HARPS-N spectra. The full table is available online, only a portion is shown for references.

BJD-2450000.0	$\log R'_{\text{HK}}$	$\sigma_{\log R'_{\text{HK}}}$	$H\alpha$	$\sigma_{H\alpha}$	HeI	$\sigma_{\text{HeI}}$	NaID	$\sigma_{\text{NaID}}$
6379.428774	−4.351	0.013	0.2462	0.0022	0.489	0.008	−	−
6380.432878	−4.273	0.023	0.2412	0.0033	0.489	0.013	0.2123	0.0050
6381.399079	−4.377	0.021	0.2369	0.0029	0.482	0.011	0.2000	0.0039
...	...	...	...	...	...	...	...	...

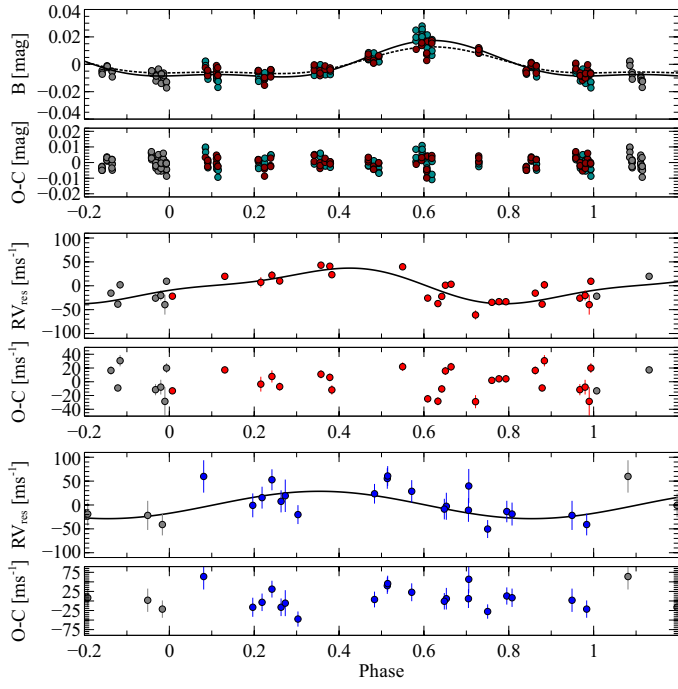


**Fig. 5.** As in Figure 4, but with the activity indicators  $\log R'_{\text{HK}}$ ,  $H\alpha$ , HeI and NaID (from top to bottom).

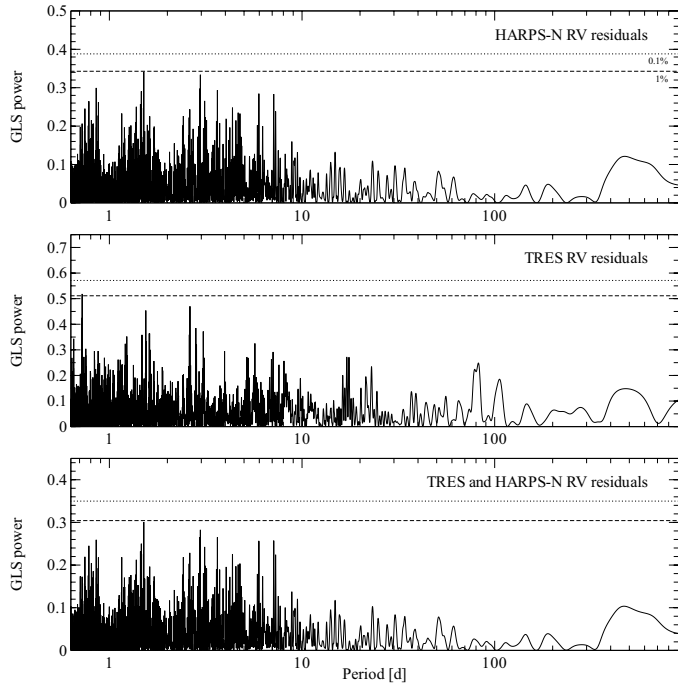


**Fig. 9.** As in Figure 8 but for the 2014 observational season.





**Fig. 10.** As in Figure 8 but for the 2015 observational season. Data in the  $r$  filter in the photometric light curve are represented by dark red points, while the activity model in this band is marked with a dotted line.



**Fig. 12.** GLS periodograms of the RV residuals after subtracting the 2-planets solution and the activity model. In the case of TRES we have included only RVs in the temporal range covered by the activity model, i. e., data gathered since 2013.